

# Power Series Solutions to Generalized Abel Integral Equations

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## Abstract

Even though they have a rather specialized structure, Abel equations form an important class of integral equations in applications. This happens because completely independent problems lead to the solution of such equations. In this paper we consider the generalized Abel integral equation of the first and second kind. Authors have been proposed a new method for constructing solutions of Abel by a power series.

## Keywords

Generalized Abel Integral Equations, Integral Equation, Power Series

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## 1. Introduction

The real world problems in scientific fields such as solid state physics, plasma physics, fluid mechanics, chemical kinetics and mathematical biology are nonlinear in general when formulated as partial differential equations or integral equations. Abel's integral equation occurs in many branches of scientific fields [1], [2] such as microscopy, seismology, radio astronomy, electron emission, atomic scattering, radar ranging, plasma diagnostics, X-ray radiography, and optical fiber evaluation. Abel's integral equation is the earliest example of an integral equation [5], [9]. In this paper, we use the method of generalized power series, to solve linear Volterra integral equations of the first and second kind. This power series are undetermined coefficients method, or a method based on the application of the Taylor series. The result obtained in the form of generalized power series solution further converted to the inversion formula of the integral equation. One such method is the representation of the solution of the equation in the form of a power series [8], [10]. Moreover, the basic theorems of this section are given without proof.

## 2. Main Results

Consider the generalized Abel integral equation of the first kind:

$$\int_a^x \frac{\phi(t)dt}{(x-t)^\alpha} = f(x), \quad (1)$$

where  $0 < \alpha < 1$  – arbitrary real constant,  $f(x): [a, b] \rightarrow R$  is given function.

Let's assume that the function  $f(x)$  can be represented as follows:

$$f(x) = c_0 + c_1(x-a) + c_2(x-a)^2 + \dots + c_n(x-a)^n + \dots \quad (2)$$

We shall seek for a solution equation in the form of the following generalized power series:

$$\phi(t) = a_0(t-a)^\gamma + a_1(t-a)^{1+\gamma} + a_2(t-a)^{2+\gamma} + \dots + a_n(t-a)^{n+\gamma} + \dots, \quad (3)$$

where  $a_n$  – the unknown coefficients that must be determined.

Substituting the power series (2), (3) into equation (1), we

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Similarly to the first case, substituting the power series (2), (7) into equation (6), performing the calculations, obtain

$$\begin{aligned} \sum_{n=0}^{+\infty} c_n x^n &= x^{\frac{1}{2}} \sum_{n=0}^{+\infty} a_n x^n + \sum_{n=0}^{+\infty} b_n x^n - \int_0^x \frac{x^{\frac{1}{2}} \sum_{n=0}^{+\infty} a_n x^n + \sum_{n=0}^{+\infty} b_n x^n}{\sqrt{x-t}} dt = \\ &= x^{\frac{1}{2}} \sum_{n=0}^{+\infty} a_n x^n + \sum_{n=0}^{+\infty} b_n x^n - \left[ \int_0^x \frac{x^{\frac{1}{2}} \sum_{n=0}^{+\infty} a_n x^n}{\sqrt{x-t}} dt + \int_0^x \frac{\sum_{n=0}^{+\infty} b_n x^n}{\sqrt{x-t}} dt \right] \Bigg|_{\substack{t=x-xz \\ dt=-xdz \\ x-t=xz \\ t=x \rightarrow z=0 \\ t=0 \rightarrow z=1}} = \\ &= x^{\frac{1}{2}} \sum_{n=0}^{+\infty} a_n x^n + \sum_{n=0}^{+\infty} b_n x^n - \sum_{n=0}^{+\infty} a_n x^{n+1} \int_0^1 z^{-\frac{1}{2}} (1-z)^{n+\frac{1}{2}} dz - \sum_{n=0}^{+\infty} b_n x^{n+\frac{1}{2}} \int_0^1 z^{-\frac{1}{2}} (1-z)^n dz = \\ &= x^{\frac{1}{2}} \sum_{n=0}^{+\infty} a_n x^n + \sum_{n=0}^{+\infty} b_n x^n - \sum_{n=0}^{+\infty} a_n x^{n+1} \mathbf{B}\left(\frac{1}{2}, n+\frac{3}{2}\right) - \sum_{n=0}^{+\infty} b_n x^{n+\frac{1}{2}} \mathbf{B}\left(\frac{1}{2}, n+1\right) \end{aligned}$$

Then by equating the terms with identical powers of  $x$  we get:

$$\begin{aligned} a_m &= b_m \mathbf{B}\left(\frac{1}{2}, m+1\right), \quad b_0 = c_0, \\ b_{m+1} &= a_m \mathbf{B}\left(\frac{1}{2}, m+\frac{3}{2}\right) + c_{m+1} = b_m \mathbf{B}\left(\frac{1}{2}, m+1\right) \mathbf{B}\left(\frac{1}{2}, m+\frac{3}{2}\right) + c_{m+1}, \\ \sum_{n=0}^{+\infty} b_n x^n &= \sum_{p=0}^{+\infty} \sum_{k=0}^p c_k x^p \pi^{p-k} \frac{k!}{p!} = \sum_{p=0}^{+\infty} \sum_{k=0}^p c_k x^p \frac{\pi^{p-k}}{(p-k-1)!} \cdot \frac{\Gamma(k+1)\Gamma(p-k)}{\Gamma(p+1)} = \\ &= \sum_{p=0}^{+\infty} \sum_{k=0}^p c_k x^p \frac{\pi^{p-k}}{(p-k-1)!} \cdot \mathbf{B}(k+1, p-k) = |n=p-k-1| = \\ &= \sum_{k=0}^{+\infty} c_k x^k + \sum_{k=0}^{+\infty} \sum_{n=0}^{+\infty} c_k x^{n+k+1} \frac{\pi^{n+1}}{n!} \mathbf{B}(k+1, n+1), \\ x^{\frac{1}{2}} \sum_{n=0}^{+\infty} a_n x^n &= x^{\frac{1}{2}} \sum_{p=0}^{+\infty} \sum_{k=0}^p c_k x^p k! \pi^{p-k} \frac{2^{p+1}}{(2p+1)!!} = \\ &= x^{\frac{1}{2}} \sum_{p=0}^{+\infty} \sum_{k=0}^p c_k x^p \pi^{p-k} \frac{\Gamma\left(\frac{1}{2}\right)\Gamma(k+1)}{\Gamma\left(k+\frac{3}{2}\right)} \cdot \frac{1}{(p-k-1)!} \cdot \frac{\Gamma\left(k+\frac{3}{2}\right)\Gamma(p-k)}{\Gamma\left(p+\frac{3}{2}\right)} = \\ &= x^{\frac{1}{2}} \sum_{p=0}^{+\infty} \sum_{k=0}^p c_k x^p \mathbf{B}\left(\frac{1}{2}, k+1\right) \cdot \frac{\pi^{p-k}}{(p-k-1)!} \cdot \mathbf{B}\left(k+\frac{3}{2}, p-k\right) = |n=p-k-1| = \\ &= \sum_{k=0}^{+\infty} c_k x^{k+\frac{1}{2}} \mathbf{B}\left(\frac{1}{2}, k+1\right) + \sum_{k=0}^{+\infty} \sum_{n=0}^{+\infty} c_k x^{n+k+\frac{3}{2}} \mathbf{B}\left(\frac{3}{2}, k+1\right) \frac{\pi^{n+1}}{n!} \mathbf{B}\left(k+\frac{3}{2}, n+1\right) \end{aligned}$$

and this implies that our solution can be written as

$$y(x) = \sum_{k=0}^{+\infty} c_k x^{k+\frac{1}{2}} \mathbf{B}\left(\frac{1}{2}, k+1\right) + \sum_{k=0}^{+\infty} \sum_{n=0}^{+\infty} c_k x^{n+k+\frac{3}{2}} \mathbf{B}\left(\frac{3}{2}, k+1\right) \frac{\pi^{n+1}}{n!} \mathbf{B}\left(k+\frac{3}{2}, n+1\right) +$$

$$\begin{aligned}
& + \sum_{k=0}^{+\infty} c_k x^k + \sum_{k=0}^{+\infty} \sum_{n=0}^{+\infty} c_k x^{n+k+1} \frac{\pi^{n+1}}{n!} B(k+1, n+1) \\
F(x) &= \sum_{k=0}^{+\infty} c_k x^k + \sum_{k=0}^{+\infty} c_k x^{k+\frac{1}{2}} B\left(\frac{1}{2}, k+1\right) = f(x) + \sum_{k=0}^{+\infty} x^{k+\frac{1}{2}} \int_0^1 c_k t^k (1-t)^{-\frac{1}{2}} dt = \\
&= f(x) + \int_0^x \frac{\sum_{k=0}^{+\infty} c_k t^k}{\sqrt{x-t}} dt = f(x) + \int_0^x \frac{f(t) dt}{\sqrt{x-t}} \\
I &= \sum_{k=0}^{+\infty} \sum_{n=0}^{+\infty} c_k x^{n+k+1} \frac{\pi^{n+1}}{n!} B(k+1, n+1) + \sum_{k=0}^{+\infty} \sum_{n=0}^{+\infty} c_k x^{n+k+\frac{3}{2}} B\left(\frac{3}{2}, k+1\right) \frac{\pi^{n+1}}{n!} B\left(k+\frac{3}{2}, n+1\right) = \\
&= \sum_{k=0}^{+\infty} \sum_{n=0}^{+\infty} c_k x^{n+k+1} \frac{\pi^{n+1}}{n!} \int_0^1 t^k (1-t)^n dt + \sum_{k=0}^{+\infty} \sum_{n=0}^{+\infty} c_k x^{n+k+\frac{3}{2}} B\left(\frac{3}{2}, k+1\right) \frac{\pi^{n+1}}{n!} \int_0^1 t^{k+\frac{1}{2}} (1-t)^n dt = \\
&= \int_0^1 \sum_{k=0}^{+\infty} \sum_{n=0}^{+\infty} c_k x^{n+k+1} \frac{\pi^{n+1} t^k (1-t)^n}{n!} dt + \int_0^x \sum_{n=0}^{+\infty} \frac{\pi^{n+1} (x-t)^n}{n!} \sum_{k=0}^{+\infty} c_k t^{k+\frac{1}{2}} B\left(\frac{3}{2}, k+1\right) dt = \\
&= \int_0^x \sum_{n=0}^{+\infty} \frac{\pi^{n+1} (x-t)^n}{n!} \sum_{k=0}^{+\infty} c_k t^k dt + \pi \int_0^x e^{\pi(x-t)} \sum_{k=0}^{+\infty} c_k t^{k+\frac{1}{2}} B\left(\frac{3}{2}, k+1\right) dt = \\
&= \pi \int_0^x e^{\pi(x-t)} [f(t) + F(t) - f(t)] dt = \pi \int_0^x e^{\pi(x-t)} F(t) dt
\end{aligned}$$

And as a final result:

$$y(x) = f(x) + \int_0^x \frac{f(t) dt}{\sqrt{x-t}} + \pi \int_0^x e^{\pi(x-t)} F(t) dt. \quad (8)$$

This solution is identical to the solution that is obtained in [4].

*Theorem 2.* If  $f(t) \in C[a, b]$ , then there exist a unique solution of equation Abel of second kind, which is expressed in the form (7).

As in first case, the theorem is proved by direct verification that the formula (8) is a solution of equation (6).

### 3. Conclusions

In this paper, solution is obtained by power series method. This may be used in more combinatorial way to obtain solution of higher degree non-linear integral equations.

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